Nuclear Astrophysics Theory - II







Image credit: Daria Sokol/MIPT Press Office

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Exotic Beam Summer School Lecture, University of Notre Dame June 7, 2022

Outline for lecture II

- Stellar burning and stellar evolution [3-9]
- Supernovae: types and nucleosynthesis [11-16]
- Making the heaviest elements: neutron capture nucleosynthesis [18-26]
- Neutron star mergers: gravitational waves, kilonovae, and nucleosynthesis [28-46]
- Galactic chemical evolution [49-51]

Stellar fusion







• The energy released by fusion provides an outward pressure, combating the gravitational inward pull



 $^{2}H + ^{3}H \rightarrow ^{4}He + n$, Q = 17.6 MeV

Energy released by fusion ~10-30 MeV



Stellar Nucleosynthesis

Evolutionary Time Scales for a 15 M_{sun} Star

Fused	Products	Time	Temperature
н	⁴ He	10^7 yrs.	4 X 10 ⁶ K
⁴ He	¹² C	Few X 10 ⁶ yrs	1 X 10 ⁸ K
¹² C	¹⁶ O, ²⁰ Ne, ²⁴ Mg, ⁴ He	1000 yrs.	6 X 10 ⁸ K
²⁰ Ne +	¹⁶ O, ²⁴ Mg	Few yrs.	1 X 10° K
¹⁶ O	²⁸ Sĩ, ³² S	One year	2 X 10 ⁹ K
²⁸ Si +	⁵⁶ Fe	Days	3 X 10 ⁹ K
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9 -⁸⁴Kr 56Fe 34S 119Sn 160 205Ti 8 12C 235U Binding energy per nucleon (MeV) Most stable nucleus 14N 238U Fusion Fission 6Li 5 Region of very stable nuclides ³He 2 ^{2}H 1-0 0 40 20 60 80 100 120 140 160 180 200 220 240 Mass number (A)

Why an iron core?



Discovering the Universe, Eighth Edition

© 2008 W. H. Freeman and Company

Neutron stars:

- Mostly neutrons, held up by neutron degeneracy pressure
- One teaspoon contains the mass of ~700 Great Pyramids
- ~10 mile diameter

C-O and O-Ne White dwarfs:

- Progenitor not able to proceed to fusion of heavier species, held up by electron degeneracy pressure
- ~200,000 times as dense as Earth with about same radius

Red dwarf:

- Most common in Solar neighborhood
- Burn H but can't reach He burning

Brown dwarf: Not able to burn H



Hertzsprung-Russell (HR) Diagrams



Hertzsprung-Russell (HR) Diagrams





Hertzsprung-Russell (HR) Diagrams



Stars in globular cluster M 3: these all formed around the same time and some have moved on to later stages, most still in hydrogen burning phase on main sequence

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Learn More

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Thermonuclear Supernovae (Type Ia)

Double degenerate model

Single-degenerate model



- Single-degenerate model:
 - C-O white dwarf accretes H- or He-rich matter from a companion (main sequence star, red giant, or helium star); mass of white dwarf increases until approaches Chandrasekhar limit $(1.4 M_{\odot})$, triggering explosion
 - Explains similar peak luminosity and early spectra for SN-Ia since 1.4 M_{\odot} implies natural limit on ${}^{56}Ni$
 - main problem is must accrete 0.3 M_{\odot} to explode since max white dwarf mass is 1.1 M_{\odot}
- Double-degenerate model:
 - Two C-O white dwarfs merge due to gravitational wave radiation, triggering explosion
 - Does not easily explain similar peak luminosity of SN-Ia due to wide range of ⁵⁶Ni production

Core-collapse SN and neutrino-driven winds



- The proto-neutron star cools through neutrino emission (99% star's binding energy released as v)
- Cooling via *Urca processes* (lepton + baryon \rightarrow baryon + v) as well as $e^+e^- \rightarrow v_l \bar{v}_l$ where $l = e, \mu, \tau$

Core-collapse SN and neutrino-driven winds





Woosley&Janka 06; see also Panov&Janka 08 Neutrinos set the neutron to proton ratio

$$Y_e = \frac{n_p}{n_p + n_n}$$

via weak interactions

$$v_e + n \rightarrow p + e^-$$

 $\bar{v}_e + p \rightarrow n + e^+$

and the influence of these reactions depends on the neutrino luminosities and average energies

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Supernovae and heavy elements? Light heavy elements and (α, n) in core-collapse SN

Conditions which synthesize A>130 are not found by most modern core-collapse SNe simulations

(e.g. Arcones+07, Wanajo+09, Fischer+10, Hüdepohl+10)

In such events other processes such as (α, n) and νp process could reach up to A~100 (e.g. Pruet+06, Fröhlich+06, Bliss+18)

Recent simulations (below) find some cases develop neutrino driven winds but not a standard feature for successful explosions



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"SWASI" analog of the SASI (Standing Accretion Shock Instability) in CCSNe – non-radial perturbations amplify leading to convective motion which impacts how the shock can move out

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Neutron capture processes to make the heaviest elements



Burbidge, Burbidge, Fowler, and Hoyle (B²FH) (1957)

Smith&Rehm 01

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Slow neutron capture (s-process) pathway



- s-process number density of neutrons ~10⁸ cm⁻³ (compare to ~10²⁴ cm⁻³ for the r process)
- Capture is "slow" relative to βdecay; implies a path close to stable species
- Note how the different paths imply some isotopes to be s-only or r-only
- s-process "seeds" are heavy nuclei such as ⁵⁶Fe (star enriched by past events)

Courtesy of Maria Lugaro

Slow neutron capture (s-process) in AGB stars

time



How do we know there is an *r*-process? Actinides (Z=89-103)

The *s*-process terminates at Pb-208 (Z=82) but *we observe actinides* in meteorites, Earth ocean crusts, our Sun, and other stars



"Curious Marie" sample of Allende meteorite shows excess U-235 which is a trace of Cm-247



Actinide boost stars compared to solar





age (Ma)

like in Reticulum II (MW in grey)

Some candidate sites for *r*-process element production

Collapsar disk winds

Magneto-rotationally driven (MHD) supernovae

Primordial black hole + neutron star







Credit: APS/Alan Stonebraker, via Physics

Siegel+18; see also McLaughlin&Surman 05, Miller+19 Winteler+12; see also Mosta+17

Fuller+17

Spotlight on MHD supernovae

Whether MHDs undergo only a "weak" *r* process reaching the second peak rather than a "main" or "strong" *r* process reaching the third peak or beyond depends on the influence of neutrinos and the magnetic field strength



Just like in CCSNe, neutrino energies and luminosities are crucial to determine the *r*-process reach



Simulations with higher magnetic field strength (ex 350C-Rs \rightarrow 10¹² G) undergo a stronger *r* process than those with lower magnetic field strength (ex 350C-Rw \rightarrow 10¹⁰ G)

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Neutron star mergers and the r process: a bit of history

M_{dyn} NS BH

Neutron-rich ejecta from neutron stars > 40 years ago

Lattimer&Schramm (1974): ~5% of the neutron star ejected as n-rich matter

Lattimer+ (1977): initially cold, expanding neutron star matter \rightarrow fission cycling *r* process capable of super heavy element formation



NSM dynamical ejecta



Effect of neutrinos



Equation of state



See also Wanajo+14, Vincent+19, Foucart+20....

Radice+19; see also Perego+19

Post-merger disk ejecta

Neutrino driven vs viscous





Miller+19

t \simeq 100 ms -100 -50 0 50 x [km] Most+21

-150

Equation of state

TNTYST

 $m t\,\simeq\,100\,ms$

 $BHB\Lambda\Phi$

GW170817 & AT2017gfo: photon opacity

Opacity sources include (*most important in NSM ejecta):

- bound-bound transitions* photoelectric absorption: photon absorbed or emitted as an electron moves between levels
- **bound-free** photoionization: electron absorbs photon and escapes
- free-free scattering bremsstrahlung: free electron passing close to ion or nucleus can emit or absorb a photon
- *electron scattering* inelastic (Compton) scattering and elastic (Rayleigh) scattering: photons scatter off electrons







GW170817 & AT2017gfo: "red" and "blue" kilonovae

Spectra and light curves depend on the species present; Lanthanide and/or actinide mass fraction \uparrow , opacity \uparrow , longer duration light curve shifted toward infrared

(e.g. Metzger+10, Lippuner+15, Barnes+16,21, Wanajo+18, Watson+19, Hotokezaka+20, Korobkin+20, Zhu+18,21, Wang+20)



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Are actinides produced in NSMs?



Are actinides produced in NSMs?



Do binary NSMs make enough heavy elements?









Impact of nuclear physics uncertainties: *r*-process N=126 peak example



- Little to no experimental data on the neutron-rich side at N=126; nuclear mass models predict different shell closure strengths and thus different amounts of elements like gold and platinum
- The N=126 shell closure is the "gateway" to the actinides and thus affects how strongly elements like uranium-238 are produced

Sensitivity of *r*-process abundances to neutron capture and β -decay



Spotlight on impact of nuclear masses

Masses determine key quantities that go into calculating capture and decay rates; for instance:

Neutron capture rates depend on

 $S_n(Z, A + 1) = M_{Z,A} + M_n - M_{Z,A+1}$

 $\beta^{\text{--}}\text{decay}$ rates depend on

$$Q_{\beta^-} = (M_{\text{parent}} - M_{\text{daughter}})c^2$$



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MCMC predictions for neutron-rich rare-earth masses



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MCMC results in *moderately neutron-rich* outflows



Neutron star merger accretion disk winds with: Hot = extended $(n,\gamma) \leftrightarrows (\gamma,n)$ equilibrium Cold = photodissociation falls out early

Vassh+21 (ApJ 907, 98)



Future experiment meets the *r*-process path



*reach of future experiment in key regions impacting the evolution of abundances (note moderately n-rich conditions used here)

> Movie by N. Vassh

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Where and when were the heavy elements we see in stars produced?



Supernovae as the *r*-process source? Galactic chemical evolution (GCE) and low metallicity stars



Supernovae as the *r*-process source? Galactic chemical evolution (GCE) and low metallicity stars



*The stars in the box seem to be more consistent with supernovae since neutron stars take time to merge Hydrodynamic mixing accounting for inhomogeneities in the interstellar medium could explain how *r*-process elements find their way to low metallicity regions

